

The background features a large, faint watermark of the University of Oxford seal. The seal is circular and contains the text 'UNIVERSITY OF OXFORD' around the perimeter. In the center, there is a shield with an open book, and the Latin motto 'DOMI MINA' is visible on the book's pages. Above the shield are two crowns, and below it is a single crown. The seal is rendered in a light beige color.

## Lecture 7: Bifurcations

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# Lecture 7 overview

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- We have considered the behaviour of autonomous systems with a single set of parameters, which were assumed to be known
- Now we will focus on how system behaviour changes depending on the values of the **constant parameters** of the system model
- Equilibrium points can change positions and character as the parameters change, leading to a **bifurcation** in the response
- This lecture will focus on categorizing bifurcations, and on providing criteria that can be used to classify them

# Local bifurcations

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- Until now our focus has been autonomous systems of the form

$$\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x})$$

- Recall (Lecture 1) that we can also consider  $\mathbf{f}$  to depend on a constant vector of parameters  $\mathbf{p} \in \mathbb{R}^p$ :

$$\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x}; \mathbf{p})$$

- This lecture considers the structural stability of solution topology in phase space near equilibrium points as a function of the vector  $\mathbf{p}$
- Here  $\mathbf{p}$  may also be called a **bifurcation vector** or **bifurcation parameter**, because the character of solution trajectories may branch (bifurcate) if the parameter values change

# 1-D bifurcations

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- The simplest systems to consider are autonomous systems with solutions on the (1-D) phase line

$$\dot{x} = f(x; p), \quad x, p \in \mathbb{R}$$

- A bifurcation occurs when the number or type of equilibrium points changes as parameter  $p$  is changed, e.g. stable to unstable
- Three types of 1-D bifurcation:
  - **saddle-node**
  - **transcritical**
  - **pitchfork**
- Bifurcations are analyzed using “normal forms” – standardized equations representing various classes of problem (not the same as linear system normal forms!)

# Saddle-node bifurcation

- The normal form of a system with a saddle-node bifurcation is

$$\dot{x} = p - x^2$$

- There are stationary points when  $0 = p - x^2 \Rightarrow x = \pm\sqrt{p}$

–  $p > 0 \Rightarrow$  two equilibria (one unstable, one stable)



–  $p = 0 \Rightarrow$  one equilibrium point (a saddle)



–  $p < 0 \Rightarrow$  no equilibrium points



- A **bifurcation diagram** shows positions and types of equilibria (vertical axis) as  $p$  varies (horizontal axis); solid lines show stable equilibria, dashed lines show unstable equilibria

# Saddle-node bifurcation diagram

- Normal form  $\dot{x} = p - x^2$

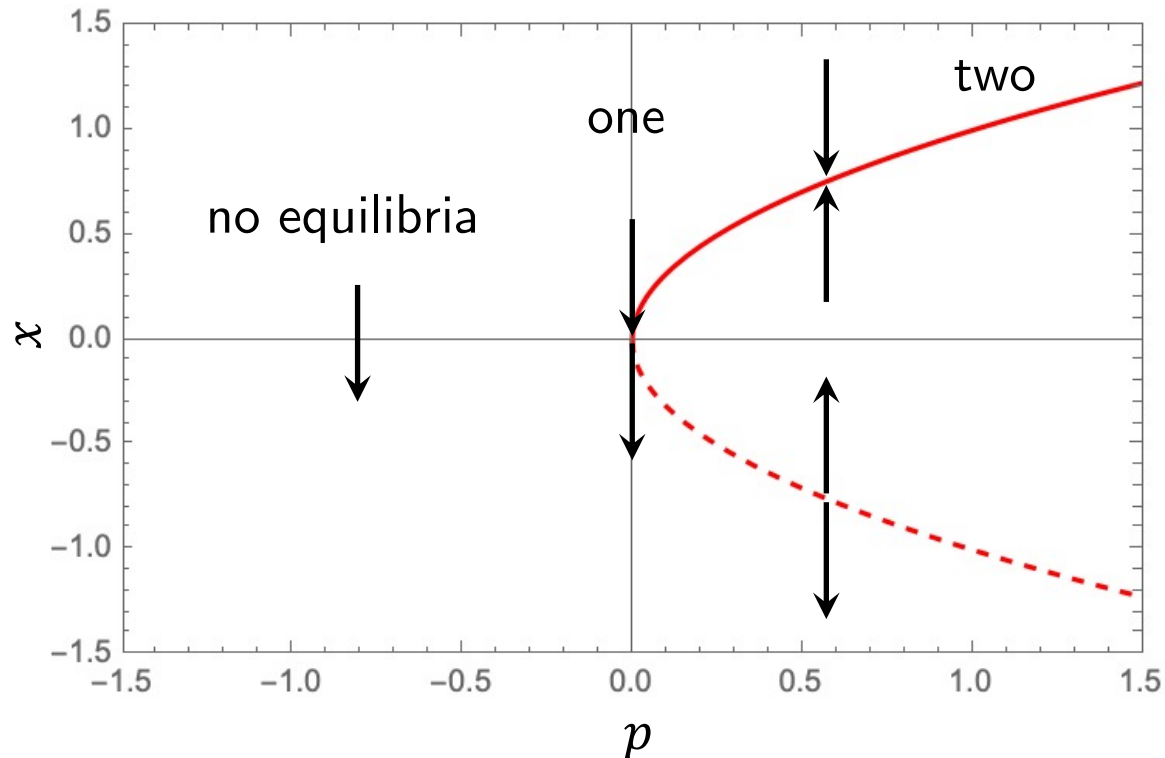
—  $p > 0$    $x$

—  $p = 0$    $x$

—  $p < 0$    $x$

characteristic flows  
on the phase line

- Bifurcation diagram



# Transcritical bifurcation

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- The normal form of a system with a transcritical bifurcation is

$$\dot{x} = px - x^2 = x(p - x)$$

- There are equilibrium points at  $x = 0$  and  $x = p$

- $p > 0 \Rightarrow$  two equilibria (one unstable, one stable)



- $p = 0 \Rightarrow$  one equilibrium point (a saddle)



- $p < 0 \Rightarrow$  two equilibria (one unstable, one stable)



- There is always a stationary point at  $x = 0$ , but its stability depends on  $p$ : the equilibria swap character as  $p$  passes through the saddle point at  $p = 0$

# Transcritical bifurcation diagram

- Normal form  $\dot{x} = px - x^2$

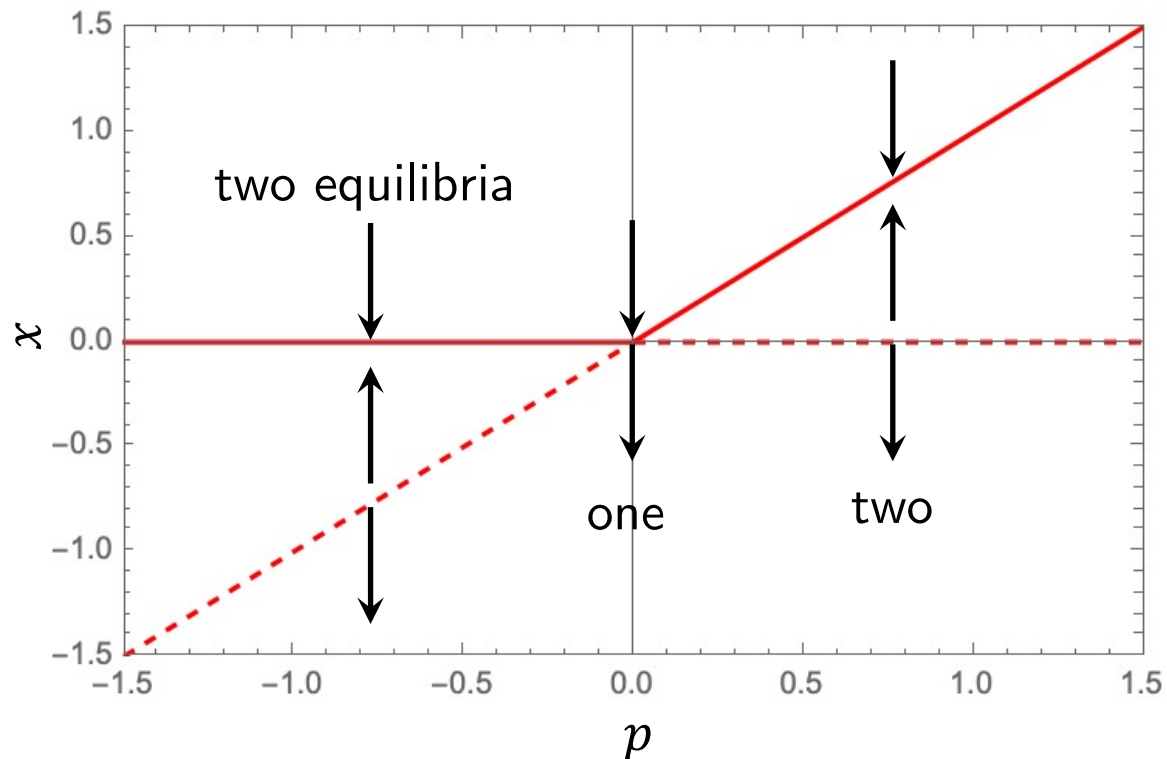
—  $p > 0$    $x$

—  $p = 0$    $x$

—  $p < 0$    $x$

characteristic flows  
on the phase line

- Bifurcation diagram





# Pitchfork bifurcation

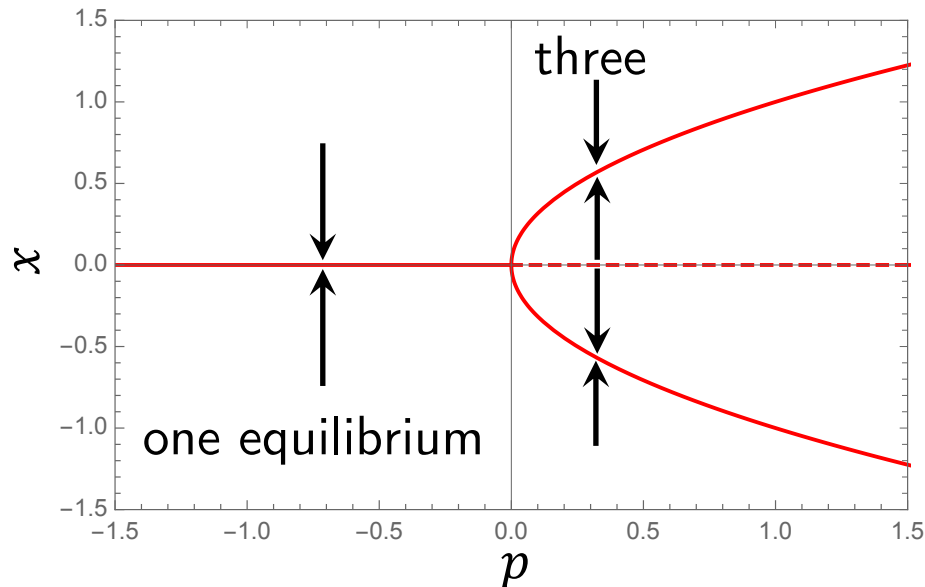
- The normal form of a system with a pitchfork bifurcation is

$$\dot{x} = px - x^3 = x(\sqrt{p} + x)(\sqrt{p} - x)$$

- There are stationary points at  $x = 0$  and, if  $p > 0$  at  $x = \pm\sqrt{p}$ 
  - $p > 0 \Rightarrow$  three equilibria (one unstable, two stable)



- $p \leq 0 \Rightarrow$  one equilibrium (stable)



# Tangency conditions

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- For a one-dimensional autonomous system, the locations  $x_0, p_0$  of bifurcation points are identified by tangency conditions

$$f(x_0, p_0) = 0 \quad \left. \frac{\partial f}{\partial x} \right|_{x_0, p_0} = 0$$

- The first condition says that  $x_0$  is an equilibrium point; the second says that  $x_0$  is a root with multiplicity two, so is non-hyperbolic

- For example consider the transcritical bifurcation:  
$$(px - x^2)|_{p_0, x_0} = x_0(p_0 - x_0) = 0$$
$$\frac{\partial}{\partial x}(px - x^2)|_{p_0, x_0} = p_0 - 2x_0 = 0 \quad \Rightarrow \quad (x_0, p_0) = (0, 0)$$

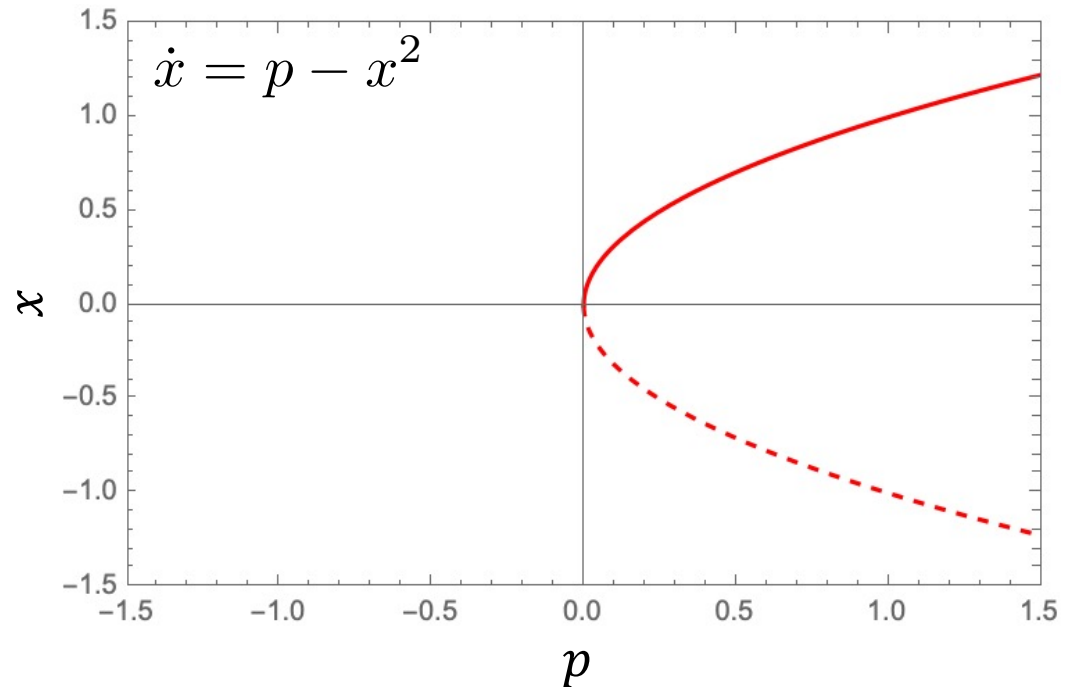
- Generally these conditions decide whether a bifurcation exists; additional conditions classify the bifurcation

# Tangency for saddle-node bifurcation

- For a bifurcation we need

$$f(x_0, p_0) = 0$$

$$\left. \frac{\partial f}{\partial x} \right|_{x_0, p_0} = 0$$



- For a saddle bifurcation we also need  $f$  to be **locally linear** in the parameter  $p$  and **locally quadratic** in the state  $x$ :

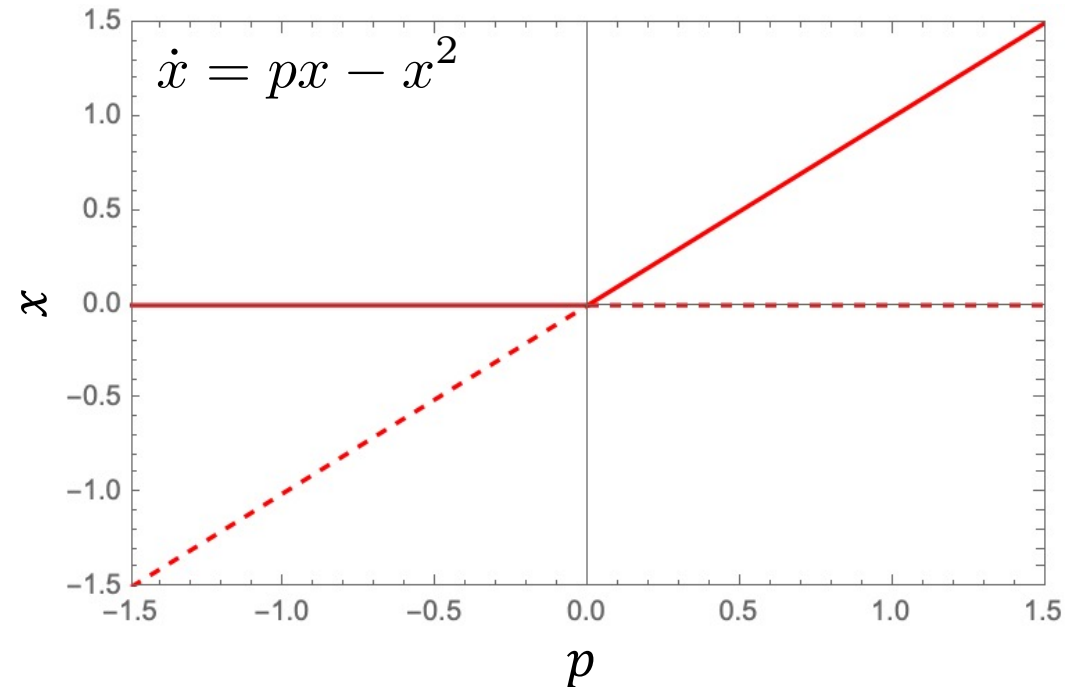
$$\left. \frac{\partial f}{\partial p} \right|_{x_0, p_0} \neq 0, \quad \left. \frac{\partial^2 f}{\partial x^2} \right|_{x_0, p_0} \neq 0$$

# Tangency for transcritical bifurcation

- For a bifurcation we need

$$f(x_0, p_0) = 0$$

$$\left. \frac{\partial f}{\partial x} \right|_{x_0, p_0} = 0$$



- For a transcritical bifurcation we also need  $f$  to be **locally bilinear** in  $x$  and  $p$ , and **locally quadratic** in the state  $x$

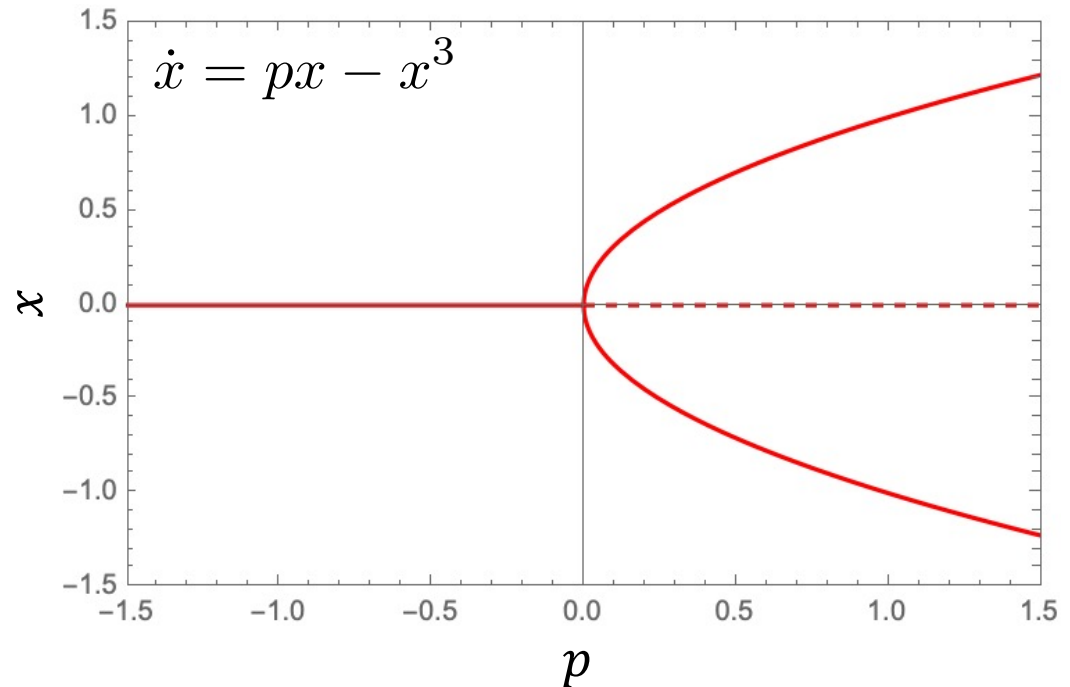
$$\left. \frac{\partial f}{\partial p} \right|_{x_0, p_0} = 0, \quad \left. \frac{\partial^2 f}{\partial x \partial p} \right|_{x_0, p_0} \neq 0, \quad \left. \frac{\partial^2 f}{\partial x^2} \right|_{x_0, p_0} \neq 0$$

# Tangency for pitchfork bifurcation

- For a bifurcation we need

$$f(x_0, p_0) = 0$$

$$\left. \frac{\partial f}{\partial x} \right|_{x_0, p_0} = 0$$

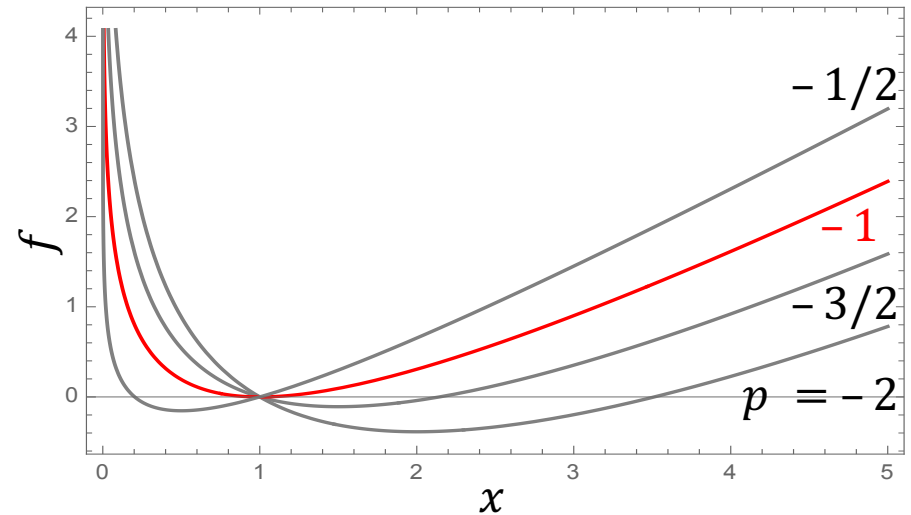


- For a pitchfork bifurcation we also need  $f$  to be **locally bilinear** in  $x$  and  $p$ , and **locally cubic** in  $x$

$$\left. \frac{\partial f}{\partial p} \right|_{x_0, p_0} = 0, \quad \left. \frac{\partial^2 f}{\partial x \partial p} \right|_{x_0, p_0} \neq 0, \quad \left. \frac{\partial^2 f}{\partial x^2} \right|_{x_0, p_0} = 0, \quad \left. \frac{\partial^3 f}{\partial x^3} \right|_{x_0, p_0} \neq 0$$

# Tangency conditions example

- Consider the system
 
$$\dot{x} = p \ln x + x - 1, \quad p < 0$$
- This has an equilibrium at  $x = 1$  and has a second equilibrium only if  $p \neq -1$



- Tangency conditions show that  $p_0 = -1$  is a **bifurcation**:

$$\begin{aligned}
 p_0 \ln x_0 + x_0 - 1 &= 0 \\
 \frac{\partial}{\partial x} (p \ln x + x) \Big|_{x_0, p_0} &= 0 \quad \implies (x_0, p_0) = (1, -1)
 \end{aligned}$$

- $$\frac{\partial}{\partial p} (p \ln x + x) \Big|_{x_0, p_0} = \ln(x_0) = 0 \quad \frac{\partial^2}{\partial x \partial p} (p \ln x + x) \Big|_{x_0, p_0} = \frac{1}{x_0} = 1$$

$$\frac{\partial^2}{\partial x^2} (p \ln x + x) \Big|_{x_0, p_0} = -\frac{p_0}{x_0^2} = 1 \quad \implies \text{transcritical}$$

# 2-D bifurcations

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- We can also characterize bifurcations for autonomous systems of higher order (than 1st order), so long as they have just a single scalar parameter  $p$
- If a 2-D system  $\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x}; p)$  has an equilibrium at  $\mathbf{x}_0$ , then bifurcations can be characterized with Sotomayor's theorem (cf. Perko 4.2), which uses the Jacobian of  $\mathbf{f}$  to formulate higher-dimensional tangency conditions
- We will not apply the theorem in detail here; instead we will explore an example of a 2-D system with a bifurcation

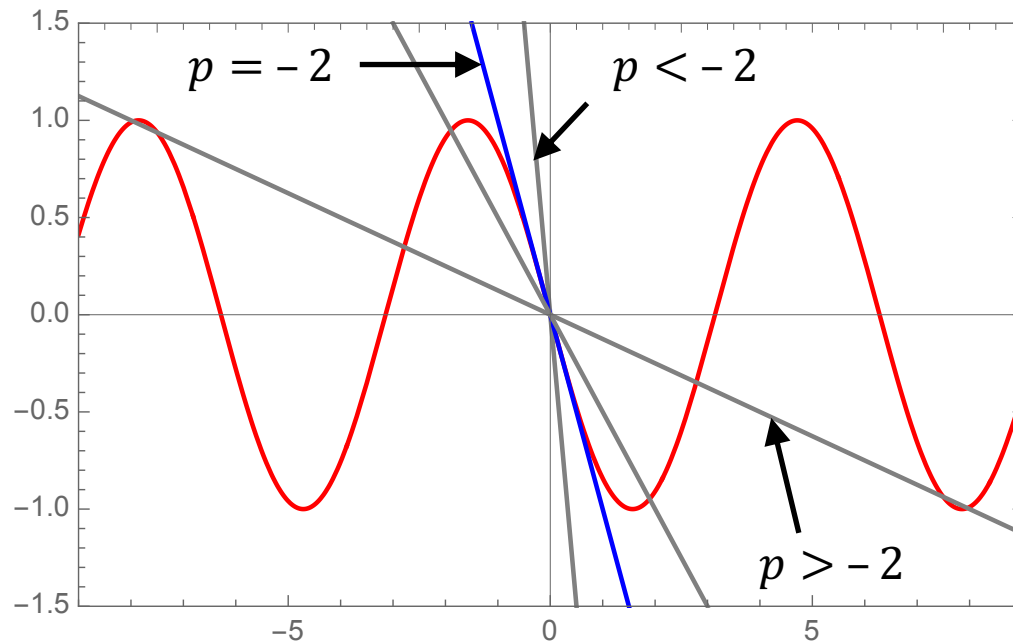
# 2-D bifurcation example

Does the origin have a bifurcation for the following system?

$$\dot{x} = px + y + \sin x$$

$$\dot{y} = x - y$$

- Find equilibria: 
$$\begin{aligned} 0 &= px_0 + y_0 + \sin x_0 \\ 0 &= x_0 - y_0 \end{aligned} \Rightarrow \begin{aligned} y_0 &= x_0 \\ (p + 1)x_0 &= -\sin x_0 \end{aligned}$$



if  $p = -2$ , the line drawn by the **left** side of this equation is tangent to the function on the **right**;  
expect bifurcation at  $p = -2$ ;



# 2-D bifurcation example

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$$\text{System } \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} px + y + \sin x \\ x - y \end{bmatrix} = \mathbf{f}(x, y)$$

- Jacobian: 
$$D\mathbf{f}(x, y) = \begin{bmatrix} p + \cos x & 1 \\ 1 & -1 \end{bmatrix}$$

$$D\mathbf{f}(0, 0) = \begin{bmatrix} p + 1 & 1 \\ 1 & -1 \end{bmatrix}$$

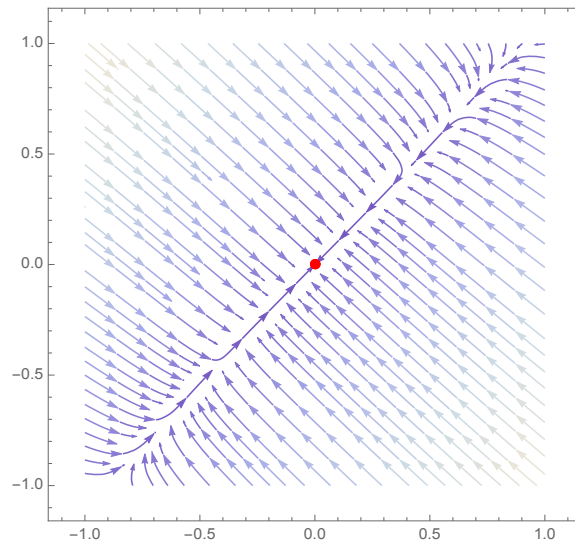
eigenvalues: 
$$\text{eig}(D\mathbf{f}(0, 0)) = \frac{1}{2}(p \pm \sqrt{(p + 2)^2 + 4})$$

- eigenvalues are: negative if  $p < -2$   
opposite in sign if  $p > -2$  (unstable direction is  $\begin{bmatrix} 1 \\ 1 \end{bmatrix}$ )
- more than one equilibrium point exists for  $p > -2$

# 2-D bifurcation example (cont'd)

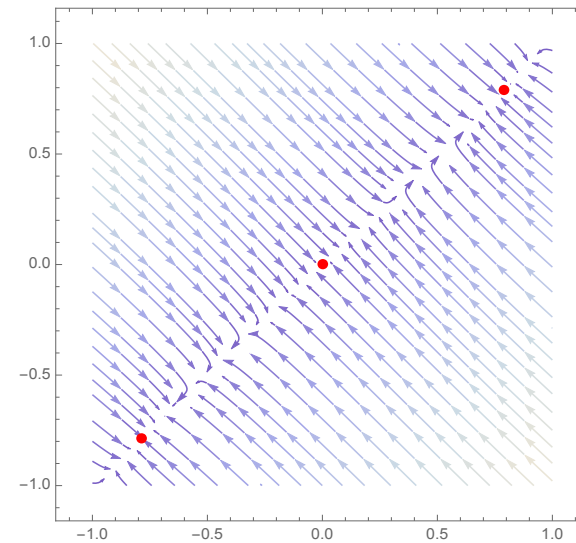
Phase portraits for  $p < -2$  and  $p > -2$ :

$p = -2.1$



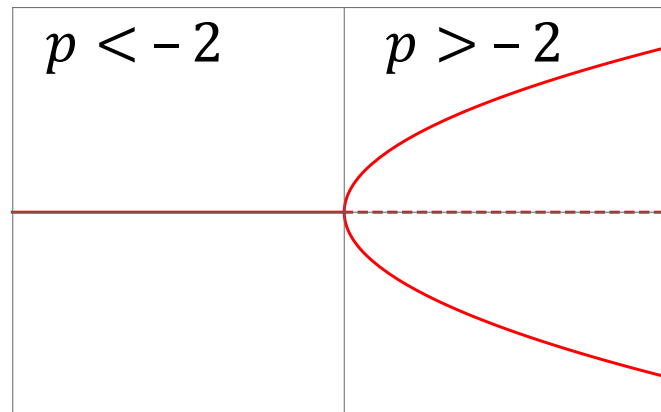
Origin is stable

$p = -1.9$



Origin is a saddle

A **pitchfork**  
bifurcation:



# Hopf bifurcations

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- The example just considered is a 2-D system with an equilibrium point at 0 that has:
  - one negative eigenvalue for all values of the parameter  $p$ ,
  - another eigenvalue passing through 0 at  $p = -2$the non-hyperbolic behaviour at  $p = -2$  was found to be a pitchfork bifurcation
- A 2-D system undergoes a **Hopf bifurcation** if the non-hyperbolic point is a centre (with pure imaginary eigenvalues)
  - in this case the stability of both eigenvalues can change

# Conditions for a Hopf bifurcation

- Assume a two-dimensional system with a scalar parameter  $p$  and equilibrium point  $\mathbf{x}^* = \mathbf{x}^*(p)$ :

$$\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x}, p), \quad \mathbf{f}(\mathbf{x}^*, p) = 0$$

- The system undergoes a Hopf bifurcation if

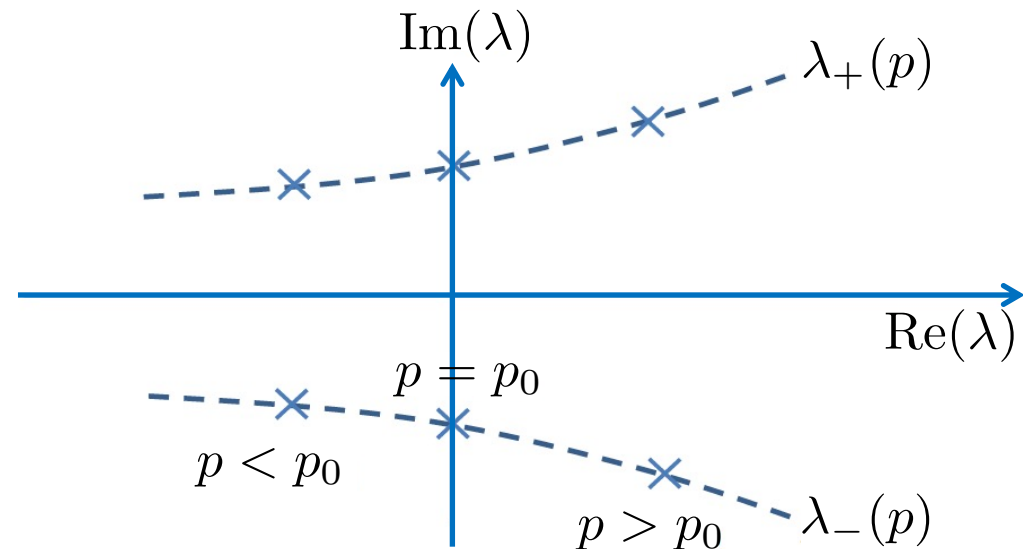
$$\text{eig}(D\mathbf{f}(\mathbf{x}^*, p)) = \lambda_{\pm}(p) = \alpha(p) \pm j\omega(p)$$

for  $p$  in the range

$$p_0 - \epsilon < p < p_0 + \epsilon$$

for some  $\epsilon > 0$  with

$$\alpha(p) \begin{cases} < 0 & \text{for } p < p_0 \\ = 0 & \text{for } p = p_0 \\ > 0 & \text{for } p > p_0 \end{cases}$$

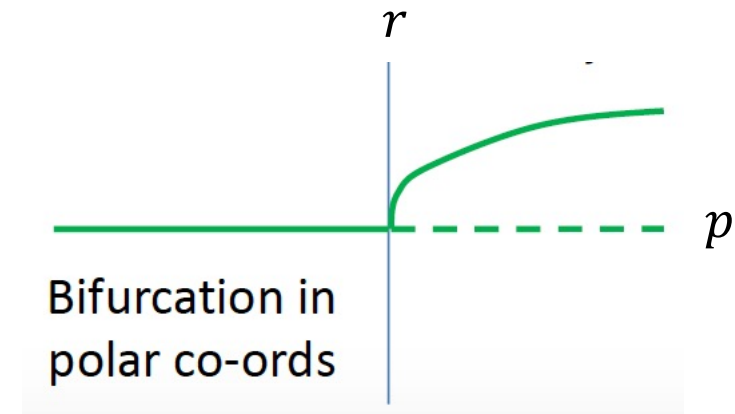
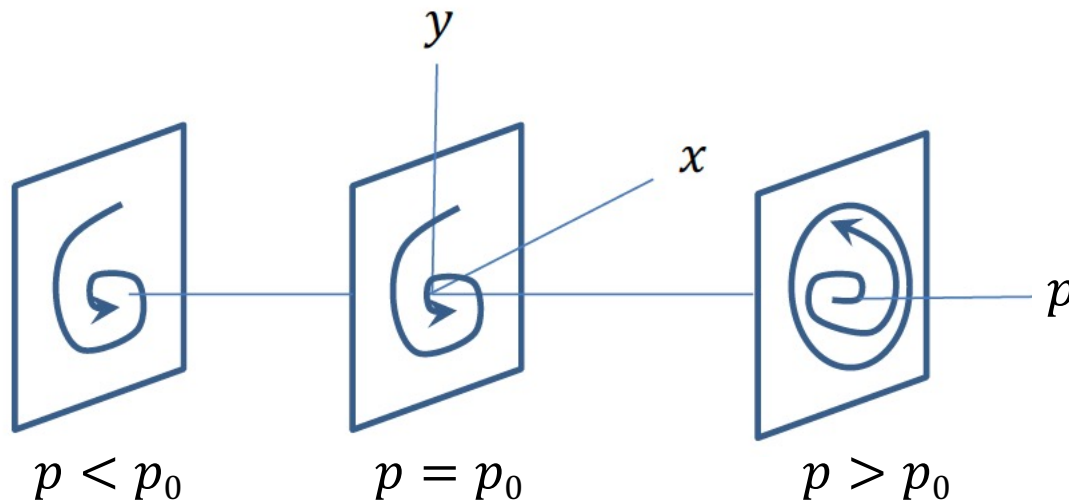


# Supercritical Hopf bifurcation

The supercritical Hopf bifurcation is best thought of in polar coordinates  $(r, \theta)$ :

- below the critical value of the parameter, there is a stable spiral equilibrium
- above the critical value, there is an unstable spiral with an enclosing stable limit cycle

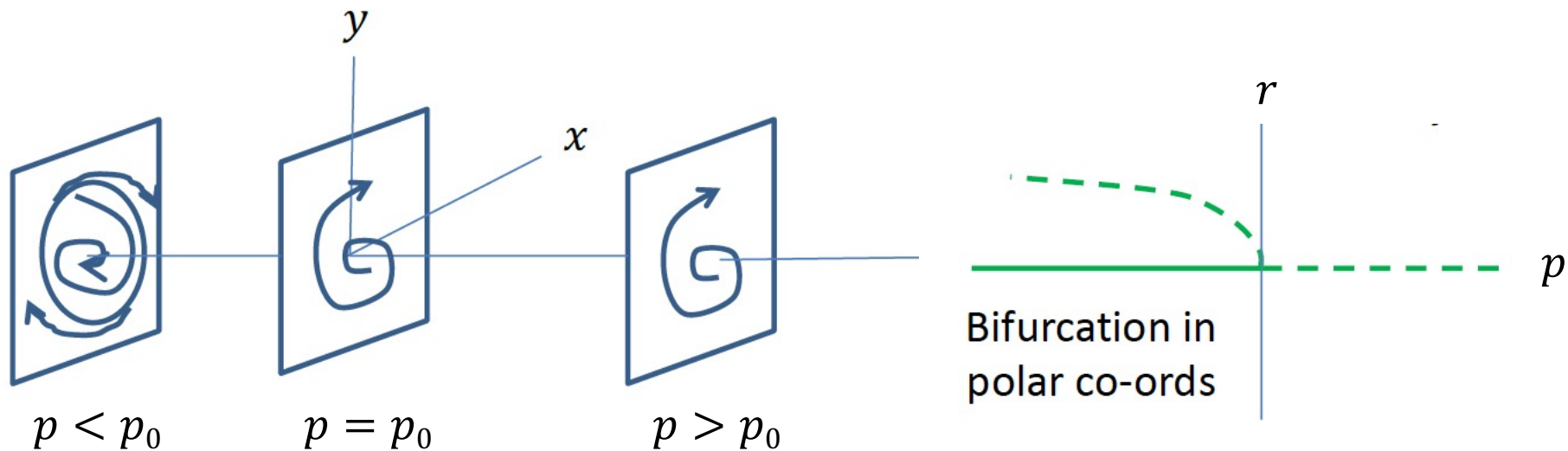
the limit cycle's radius  $r$  expands with increasing  $p$



# Subcritical Hopf bifurcation

The subcritical Hopf bifurcation behaves as follows:

- below the critical value of  $p$  there is a stable spiral surrounded by an unstable limit cycle
- the limit cycle radius shrinks as  $p$  increases
- at the critical value the cycle collapses to a fixed point
- above the critical value there is an unstable spiral



# Degenerate Hopf bifurcation

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The degenerate Hopf bifurcation behaves as follows:

- below the critical value of  $p$  there is a stable spiral
- at the critical value of  $p$  the spiral becomes a nonlinear centre whose orbit is not isolated ( $r(t)$  depends on initial conditions)
- above the critical value of  $p$  there is an unstable spiral

- Called a 'degenerate' bifurcation because there is a non-isolated orbit at the critical parameter value
- The degenerate Hopf bifurcation has no limit cycles for any value of the parameter  $p$

# Hopf bifurcation example

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Consider the system  $\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} px - y + xy^2 \\ x + py + y^3 \end{bmatrix} = \mathbf{f}(x, y)$

- Just one equilibrium point:  $\mathbf{f}(x, y) = (0, 0) \implies (x, y) = (0, 0)$
- Eigenvalues of Jacobian at  $(x, y) = (0, 0)$ :

$$D\mathbf{f}(0, 0) = \begin{bmatrix} p + y^2 & -1 + 2xy \\ 1 & p + 3y^2 \end{bmatrix} \Big|_{(0,0)} = \begin{bmatrix} p & -1 \\ 1 & p \end{bmatrix} \implies \lambda_{\pm} = p \pm j$$

- From this we expect a Hopf bifurcation at  $p = 0$



# Hopf bifurcation example (cont'd)

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What type of Hopf bifurcation does this system have?

Transform into polar coordinates:  $\dot{r} = \frac{x\dot{x} + y\dot{y}}{r} = r(p + r^2 \cos^2 \theta)$

$$\dot{\theta} = \frac{x\dot{y} - y\dot{x}}{r^2} = 1$$

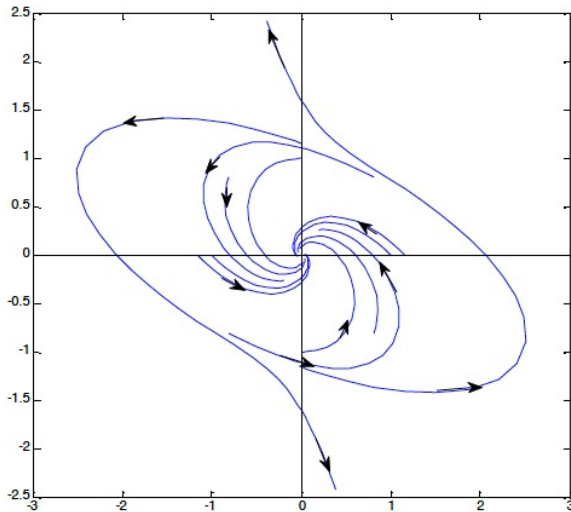
- $p > 0 \implies \dot{r} \geq pr$        $\dot{r} > 0$  for all  $t$   
so no limit cycle
- $p = 0 \implies \dot{r} \geq 0$        $\dot{r} \geq 0$  for all  $t$   
so no limit cycle
- $p < 0 \implies \dot{r} = pr + ry^2$        $\dot{r} < 0$  for  $y < |p|^{1/2}$   
so a stable spiral

Therefore a **subcritical Hopf bifurcation** occurs at  $p = 0$ ,  
so expect an unstable limit cycle

# Hopf bifurcation example: phase portraits

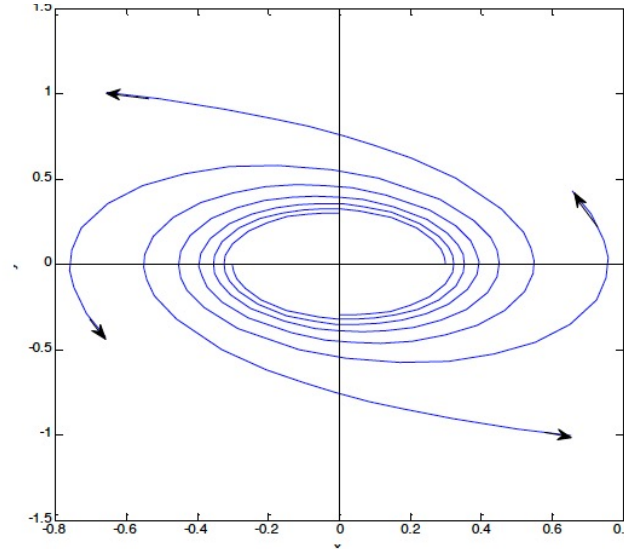
$$\begin{aligned} \dot{x} &= px - y + xy^2 \\ \dot{y} &= x + py + y^3 \end{aligned} \iff \begin{aligned} \dot{r} &= r(p + r^2 \cos^2 \theta) \\ \dot{\theta} &= 1 \end{aligned}$$

$p = -1$



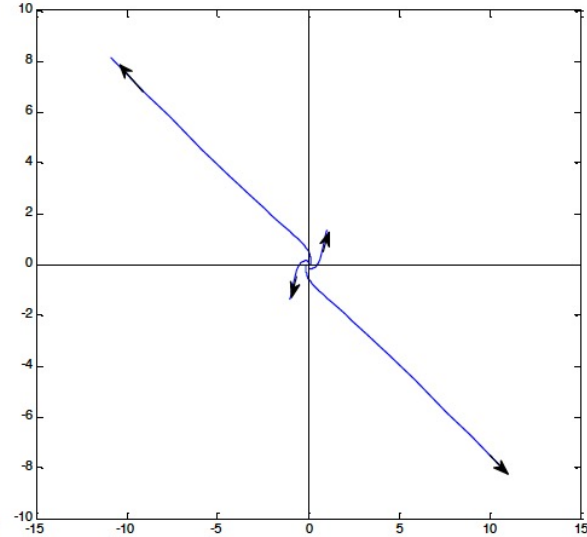
Stable spiral  
with unstable  
limit cycle

$p = 0$



Unstable spiral

$p = 1$



Unstable spiral